

Minimal Technicolor on the lattice

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Background:

- **Higgs field is special:**
 - ▶ it has not been found
 - ▶ it is scalar
 - ▶ ... hierarchy problem, vacuum stability, unitarity bound ...
- **Can we live without a scalar?**
- Consider the EW symmetry breaking and χ SB in QCD:

	EWSB	χ SB
condensate:	Higgs vev v	$\bar{\psi}\psi$ chiral condensate
goldstones:	eaten by W,Z (gauged)	π -mesons
radial excitation:	Higgs particle	scalar meson

Technicolor

Technicolor (TC): Electroweak symmetry breaking \longrightarrow chiral symmetry breaking of “Techni-QCD”, (technigauge + techniquarks Q), with $\Lambda_{TC} \approx \Lambda_{EW}$.

- Quarks have technicolor and EW charge
- After chiral symmetry breaking:
 - \Rightarrow scalar $\bar{Q}Q$ -meson: Higgs
 - \Rightarrow Pseudoscalars \rightarrow W,Z -longitudinal modes
 - \Rightarrow exotic technihadrons

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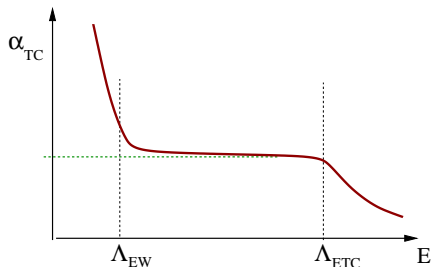
It's both Strong and Electroweak

What about Yukawa couplings to SM quarks?

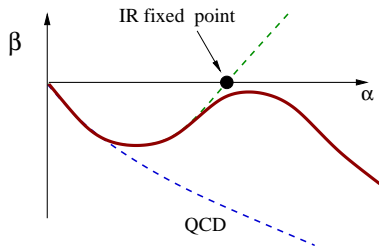
- ⇒ **Extended technicolor (ETC)**: new gauge interaction coupling normal quarks and techniquarks.
- At low energy, $\mathcal{L}_{\text{Yukawa}} \sim \langle \bar{Q}Q \rangle \bar{q}q$
 - **Experimental constraint:** $\Lambda_{\text{ETC}} \gg \Lambda_{\text{EW}} = \Lambda_{\text{TC}}$, due to FCNC's.
 - Typically $\Lambda_{\text{ETC}} \sim 100\text{--}1000 \times \Lambda_{\text{EW}}$

Walking coupling

To make this work, the theory has to be (almost) conformal in a wide range of energy:



β -function is almost zero at moderately strong coupling

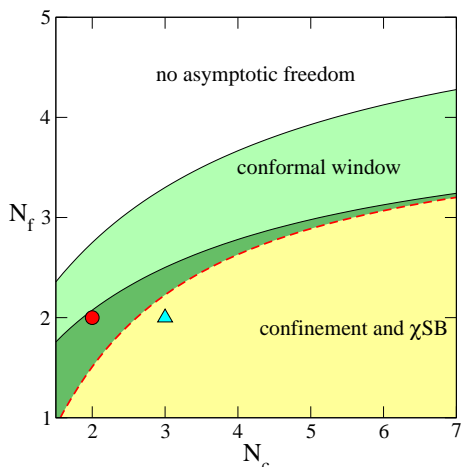
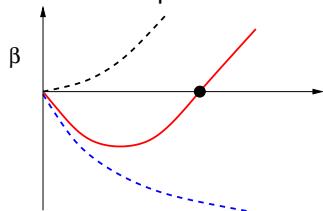


Symmetric representation

- The required behaviour is difficult to satisfy using fundamental rep. quarks. (Need large $N_f \rightarrow$ difficult to avoid FCNC's.)
- However, for $SU(N)$ with 2-index symmetric representation quarks ($\square\square$) $N_f \leq 5$ is sufficient to reach conformal behaviour for any N_c (at least in perturbative analysis).
- There \exists “conformal window” with IR fixed point
[Dietrich, Sannino, Tuominen; Dietrich, Sannino]

Conformal window

Within the conformal window
there \exists IR fixed point



Lattice simulations have been done at $N_c = 2$, $N_f = 2$ and $N_c = 3$,
 $N_f = 2$

Why lattice simulations?

- χ SB is essentially non-perturbative: lattice simulations are needed to check whether the scenario works.
- We study here $N_c = 2$, $N_f = 2$ -case; the simplest model in this class: “Minimal technicolor”.
- Also studied by [Catterall, Sannino; Del Debbio, Patella, Pica].
- $N_c = 3$, $N_f = 2$ has been studied by [DeGrand, Shamir, Svetitsky].

How to observe the “walking” /conformal behaviour?

- Measure the running coupling directly using “Schrödinger functional” (not yet done!)
- Particle spectrum: all modes become massless as $m_q \rightarrow 0$.

On the lattice:

- SU(2) Wilson gauge action in fundamental rep.
- SU(2) Wilson fermions in symmetric (\equiv adjoint for SU(2)) rep.
- volumes $20^4 - 32^4$, 5 different lattice couplings $\beta \equiv 4/g^2$.

Phase diagram

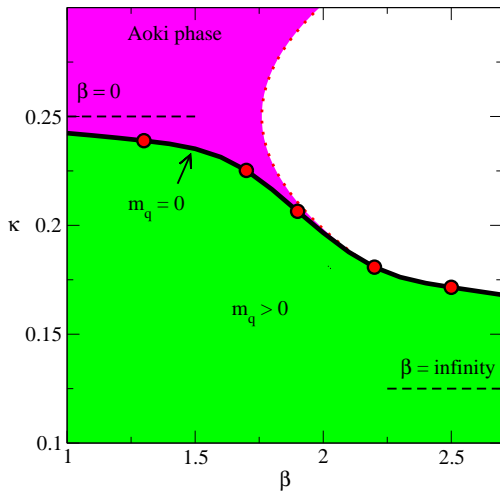
Lattice parameters:

$$\beta = \frac{4}{g^2} \quad \kappa = \frac{1}{8 + 2m_{q,\text{bare}}}$$

We determine physical quark mass from the axial Ward identity

$$m_q = \lim_{t \rightarrow \infty} \frac{1}{2} \frac{\partial_t V_{\text{PS}}(t)}{V_{\text{PP}}(t)}$$

$m_q(\beta, \kappa) = 0$ determines the critical line $\kappa_c(\beta)$



States

$SU(2)$ + fundamental quarks:

- 2-quark states $\bar{q}q$, qq (degenerate except in isoscalar channel)
- glueballs

$SU(2)$ + adjoint quarks:

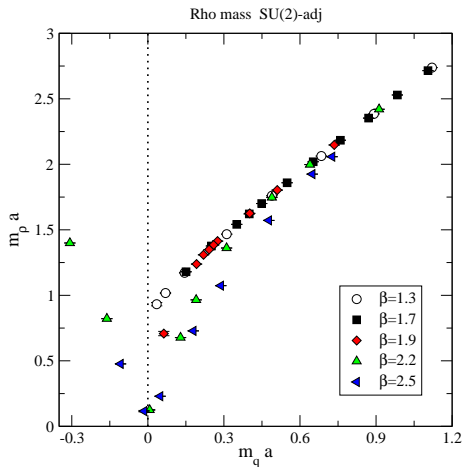
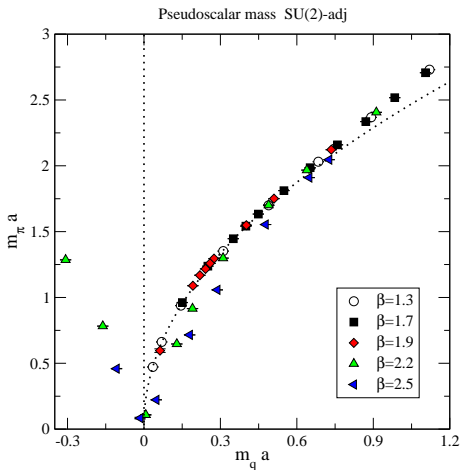
- QQ “mesons” – π , ρ . . .
- QQQ “baryons” – “proton”
- Qg quark-gluon state
- glueballs

What to expect:

- If QCD-like χ SB: as $m_Q a \rightarrow 0$,
 - ▶ $m_\pi \propto m_Q^{1/2}$
 - ▶ other states have finite mass.
- If IR fixed conformal point: $m_Q a \rightarrow 0$, all states become massless.
~free quarks?
- If walking behaviour
 - ▶ There $\exists \beta_{\text{conf.}}$, where system appears almost conformal.
 - ▶ When $\beta < \beta_{\text{conf.}}$, there is χ SB as in QCD.
 - ▶ When $\beta > \beta_{\text{conf.}}$, theory looks non- χ SB except in humongous volumes.
- On the lattice extrapolation $m_Q a \rightarrow 0$ is required. Too large m_Q or too small V can lead to misleading results.

Results

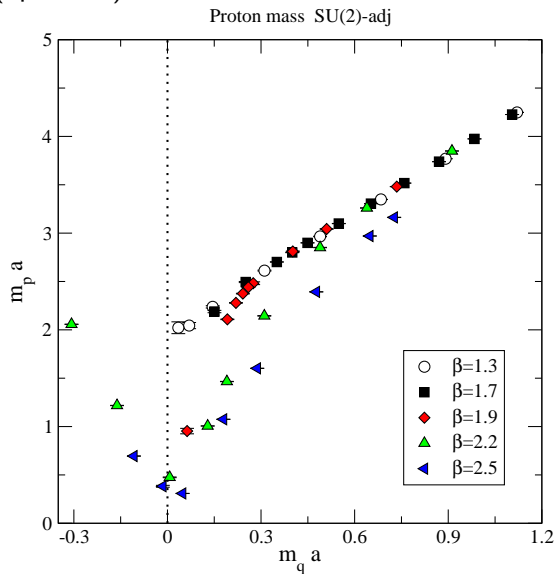
Pseudoscalar (“ π ”) and vector (“ ρ ”) masses



At small β , looks like χ SB. At large β ?

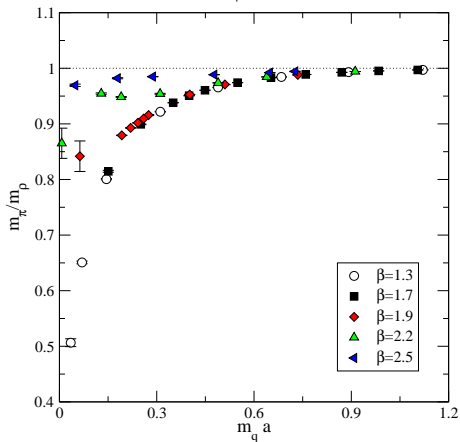
Results

3-quark state ("proton") mass

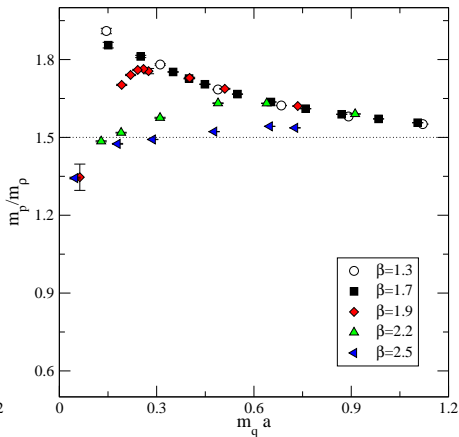


Results

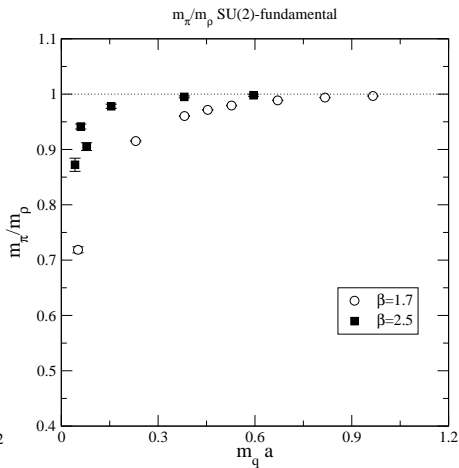
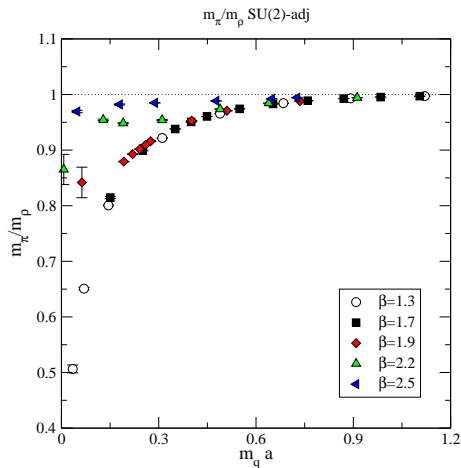
Mass ratios m_π/m_ρ and m_{Proton}/m_ρ
 m_π/m_ρ SU(2)-adj



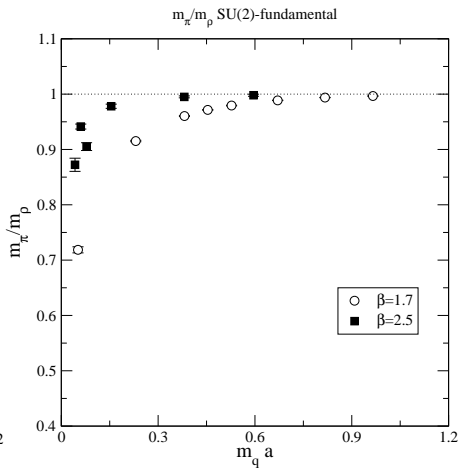
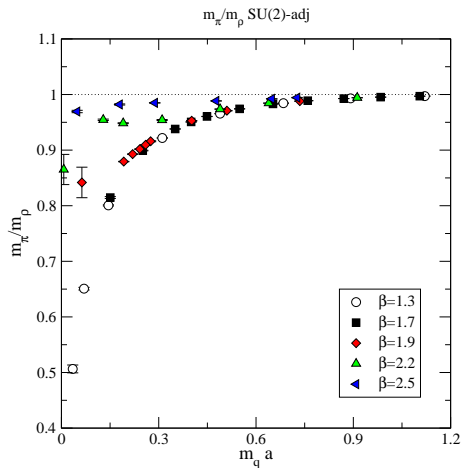
m_p/m_ρ SU(2)-adj



Compare with fundamental rep.



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Conclusions (preliminary!)

- There is χ SB at small β
 - Theory appears to be more “walking” (i.e. slower) than fundamental quark theory, but how much?
 - It may be possible that there exists almost-fixed-point near $\beta \sim 2 \Rightarrow g_{\text{Lat.}}^2 \sim 2?$
 - There is also room for a genuine IR fixed point. In that case small- β χ SB phase is a lattice artifact, and is separated by a phase transition at some β from the continuum (large β) behaviour.
 - Direct evaluation of β -function required (in progress).
-
- Predictions obtainable from this TC model:
 - ▶ Measure $\langle \bar{Q}Q \rangle$, set to v_{Higgs}
 - ▶ Measure QQ scalar mass \rightarrow Higgs mass
 - ▶ ρ : lightest exotic particle
 - ▶ Modified by ETC corrections