

**Comment on the paper “Macdonald polynomials and algebraic integrability”
by O. A. Chalykh**

Note by Tom H. Koornwinder, T.H.Koornwinder@uva.nl, December 6, 2010

Chalykh computed in [1, §4.1] a Baker-Akhiezer (BA) function for root system A_1 . It is given by his formulas (4.1), (4.4) and (4.6). Let me write his resulting BA function $\psi = \psi(x, z; m)$, depending on a parameter $m \in \mathbb{Z}_{\geq 0}$, as a ${}_2\phi_1$:

$$\psi(x, z; m) = (-1)^m q^{-\frac{3}{2}m^2 - \frac{1}{2}m} q^{(x-m)(z-m)} (q^{2x+2}; q^2)_m {}_2\phi_1 \left(\begin{matrix} q^{-2m}, q^{-2m+2x} \\ q^{2x+2} \end{matrix}; q^2, q^{2(m+z+1)} \right). \quad (1)$$

Observe the following properties of ψ .

- We have

$$\psi(x, z; m) = \psi(-x, -z; m). \quad (2)$$

Indeed, for generic x the ${}_2\phi_1$ in (1) can be written as

$$\sum_{j=0}^m \frac{(q^{-2m}; q^2)_j (q^{-2m+2x}; q^2)_j}{(q^{2x+2}; q^2)_j (q^2; q^2)_j} (q^{2(m+z+1)})^j = \sum_{j=0}^m c_j = \sum_{j=0}^m c_{m-j},$$

and then the right-hand side can again be written as a ${}_2\phi_1$, see [3, (1.8)] (or more generally [2, Exercise 1.4(ii)]). Then the resulting identity can be rewritten as (2).

- We have

$$\psi(x, z; m) = \psi(x, -z; m) \quad (x = 1, \dots, m). \quad (3)$$

Indeed, now write the ${}_2\phi_1$ in (1) as

$$\sum_{j=0}^{m-x} c_j = \sum_{j=0}^{m-x} c_{m-x-j},$$

and write the right-hand side again as a ${}_2\phi_1$. Then the resulting identity can be rewritten as (3).

- We have

$$\psi(x, z; m) = \psi(z, x; m). \quad (4)$$

Indeed, (4) is a rewritten version of the transformation formula [2, (III.2)].

- We have

$$\psi(x, z; m) = \psi(x, -z; m) \quad (z = 1, \dots, m). \quad (5)$$

Indeed, combine (3), (2) and (5).

- Note that we can use (1) as a definition of $\psi(x, z; m)$ for general integer m by the convention

$$(q^{2x+2}; q^2)_m = \frac{(q^{2x+2}; q^2)_\infty}{(q^{2x+2m+2}; q^2)_\infty}.$$

We have

$$\psi(x, z; -m) = - \frac{q^{(2m-1)(x+z)}}{(q^{2x-2m+2}; q^2)_{2m-1} (q^{2z-2m+2}; q^2)_{2m-1}} \psi(x, z; m-1) \quad (m \in \mathbb{Z}_{>0}). \quad (6)$$

Indeed, (6) is a rewritten version of [2, (III.3)]. Note that $\psi(x, z; -m)$ ($m \in \mathbb{Z}_{>0}$) has poles for $x \in \{1, 2, \dots, m-1\}$ and for $z \in \{1, 2, \dots, m-1\}$, by which there are no immediate analogues of (3) and (5) for $\psi(x, z; -m)$. However, (2) and (4) remain valid for negative integer m .

With ψ given by (1) we can now review [1, Proposition 4.2, formula (4.10) and Lemma 5.4]:

1. ψ is a normalized BA function. Indeed, (1) has the form [1, (4.1)]:

$$\psi(x, z; m) = q^{xz} \sum_{j=0}^m \psi_{-m+2j}(x; m) q^{(-m+2j)z},$$

and it satisfies the normalization condition [1, (4.10)]:

$$\psi_m(x; m) = \prod_{j=1}^m (q^{j-x} - q^{-j+x}).$$

The function ψ also satisfies condition [1, (4.2)], since this condition can be rewritten as (5).

2. ψ is symmetric in x and z , see (4).
3. ψ satisfies the difference equation

$$\frac{q^{x-m} - q^{-x+m}}{q^x - q^{-x}} \psi(x+1, z; m) + \frac{q^{x+m} - q^{-x-m}}{q^x - q^{-x}} \psi(x-1, z; m) = (q^z + q^{-z}) \psi(x, z; m). \quad (7)$$

Indeed, (7) combined with (4) is a rewritten version of [2, Exercise 1.13], see also [3, p.15 below].

4. ψ satisfies [1, Lemma 5.4]:

$$\psi(wx, wz; m) = \psi(x, z; m) \quad (w \in W).$$

Indeed, this turns down to (2).

Next I will discuss [1, Theorem 5.11] for root system A_1 . First we identify Macdonald polynomials for that root system with continuous q -ultraspherical polynomials. Following (2.10), (2.11), (2.1) and (2.13) in [1] the Macdonald polynomial then has the form

$$P_n(x; q, t) = \sum_{j=0}^n f_{n-2j} q^{(n-2j)x},$$

where $f_{n-2j} = f_{2j-n}$ and $f_n = 1$, and it satisfies

$$D P_n = (q^n t + q^{-n} t^{-1}) P_n,$$

where D is the difference operator given by

$$(Df)(x) := \frac{tq^x - t^{-1}q^{-x}}{q^x - q^{-x}} f(x+1) + \frac{tq^{-x} - t^{-1}q^x}{q^x - q^{-x}} f(x-1).$$

By the reasoning in [4, (9.10) and following] we obtain

$$P_n(x; q, t) = \frac{(q^2; q^2)_n}{(t^2; q^2)_n} C_n\left(\frac{1}{2}(q^x + q^{-x}); t^2 \mid q^2\right), \quad (8)$$

where C_n is a continuous q -ultraspherical polynomial, see [2, (7.4.2)]:

$$C_n(\cos \theta; t \mid q) := \sum_{j=0}^n \frac{(t; q)_j (t; q)_{n-j}}{(q; q)_j (q; q)_{n-j}} e^{i(n-2j)\theta}.$$

Now observe [2, (7.4.5)]:

$$C_n(\cos \theta; t \mid q) = \frac{(t; q)_\infty}{(t^2; q)_\infty} \frac{(t^2; q)_n}{(q; q)_n} \left(\tilde{D}_n(e^{i\theta}; t \mid q) + \tilde{D}_n(e^{-i\theta}; t \mid q) \right), \quad (9)$$

where

$$\tilde{D}_n(e^{i\theta}; t \mid q) := e^{in\theta} \frac{(te^{-2i\theta}; q)_\infty}{(e^{-2i\theta}; q)_\infty} {}_2\phi_1 \left(\begin{matrix} t, te^{2i\theta} \\ qe^{2i\theta} \end{matrix}; q, q^{n+1} \right). \quad (10)$$

(Note that, by [2, (III.1)], the ${}_2\phi_1$ in (10) coincides with the ${}_2\phi_1$ for $S_n(e^{i\theta}; t^{1/2}, (qt)^{1/2}, -t^{1/2}, -(qt)^{1/2})$ in [5, (3.4)].) Observe that

$$\tilde{D}_{n+m}(q^x; q^{-2m} \mid q^2) = q^{mn} q^{-\frac{1}{2}m(m+1)} \psi(x, n; m). \quad (11)$$

After replacing in (9) $q, t, e^{i\theta}, n$ by $q^2, q^{-2m}, q^x, n+m$, respectively, we arrive at the first identity in [1, Theorem 5.11] for A_1 :

$$\psi(x, n; m) + \psi(-x, n; m) = \prod_{j=1}^m (q^{j-n} - q^{-j+n}) P_{n+m}(x; q, q^{-m}). \quad (12)$$

For the derivation of the second identity in [1, Theorem 5.11] for A_1 we apply [2, (III.3)] to (10). Then we obtain in particular the identity

$$\tilde{D}_{n-m-1}(q^x; q^{2m+2} | q^2) = \frac{1}{(q^{2n-2m}; q^2)_{2m+1}} \frac{\tilde{D}_{n+m}(q^x; q^{-2m} | q^2)}{\prod_{j=-m}^m (q^{j-x} - q^{-j+x})}.$$

Substitution in (9) yields

$$\begin{aligned} (q^2; q^2)_m \prod_{j=-m}^m (q^{j-x} - q^{-j+x}) C_{n-m-1}(\tfrac{1}{2}(q^x + q^{-x}); q^{2m+2} | q^2) \\ = \tilde{D}_{n+m}(q^x; q^{-2m} | q^2) - \tilde{D}_{n+m}(q^{-x}; q^{-2m} | q^2). \end{aligned}$$

By substitution of (8) and (11) we arrive at [1, Theorem 5.11] for A_1 :

$$\psi(x, n; m) - \psi(x, -n; m) = \prod_{j=1}^m (q^{j-n} - q^{-j+n}) \prod_{j=-m}^m (q^{j-x} - q^{-j+x}) P_{n-m-1}(x; q, q^{m+1}). \quad (13)$$

References

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