## Forward is backward for time-homogeneous diffusions

 orWhy worry about boundary conditions?

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In physics, spatial boundary conditions for the heat equation on a half-line (half-space) are always provided.

In finance such boundary conditions are often omitted for pricing equations, perhaps since for geometric Brownian motion such conditions are redundant.

For many models in finance, however, boundary conditions are needed and even if they are redundant, boundary conditions (boundary behaviour) provides useful information.

We will here not focus on when the conditions are redundant or not, rather on the information they provide. This is particularly striking in the case when we consider forward equations whose boundary behaviour can be reduced, in the case of time-homogeneous diffusions, to the study of boundary behaviour for backward equations.

## Feynman-Kac

We study diffusions on the half-line satisfying

$$
d X_{t}=\beta\left(X_{t}\right) d t+\sigma\left(X_{t}\right) d W
$$

Consider an expected value of the form

$$
u(x, t)=E_{x, t} g\left(X_{T}\right)
$$

The corresponding partial differential equation is

$$
\left\{\begin{array}{l}
u_{t}+\frac{1}{2} \sigma^{2}(x) u_{x x}+\beta(x) u_{x}=0 \\
u(x, T)=g(x)
\end{array}\right.
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What about boundary conditions at $x=0$ ?

## An example: The Black-Scholes model

Stock price:

$$
d X_{t}=r X_{t} d t+\sigma X_{t} d W
$$

Call price:

$$
u(x, t)=E_{x, t} e^{-r(T-t)}\left(X_{T}-K\right)^{+}
$$

The BS-equation:

$$
\left\{\begin{array}{l}
u_{t}+\frac{1}{2} \sigma^{2} x^{2} u_{x x}+r x u_{x}-r u=0 \\
u(x, T)=(x-K)^{+}
\end{array}\right.
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Boundary condition at $x=0$ :

$$
u(0, t)=0
$$

## Another example: The CIR-process

Assume that

$$
d X_{t}=\left(b-a X_{t}\right) d t+\sigma \sqrt{X_{t}} d W
$$

and

$$
u(x, t)=E_{x, t} g\left(X_{T}\right)
$$

The corresponding partial differential equation is

$$
\left\{\begin{array}{l}
u_{t}+\frac{1}{2} \sigma^{2}(x) u_{x x}+\beta(x) u_{x}=0 \\
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$$

What about boundary conditions at $x=0$ ?

## The CIR-process, continued

- If $b<\sigma^{2} / 2$, then 0 is reached with positive probability. Boundary conditions are needed to ensure uniqueness of solutions. Indeed, Heston-Loewenstein-Willard (2007) provide two different continuous bounded solutions to the term structure equation (with no specified boundary conditions).
- If $b>\sigma^{2} / 2$, then 0 cannot be reached. Boundary conditions are not needed from a mathematical perspective, but crucial from a numerical perspective!


## Boundary conditions for the backward equation

Hypothesis
Assume that $\sigma^{2}(x, t)$ and $\beta(x, t)$ are $C^{1}$ in $x, \sigma(x, t)=0$ iff $x=0, \beta(0, t) \geq 0$ and linear growth conditions. Let $g$ be $C^{1}$ with $g$ and $g^{\prime}$ bounded.

## Theorem

The function $u(x, t)=E_{x, t} g(X(T))$ is the unique classical solution to the backward equation

$$
\left\{\begin{array}{l}
u_{t}+\frac{1}{2} \sigma^{2} u_{x x}+\beta u_{x}=0 \\
u(x, T)=g(x) \\
u_{t}(0, t)+\beta(0, t) u_{x}(0, t)=0 .
\end{array}\right.
$$

Note: A classical solution is understood to be $C^{1}$ up to the boundary $x=0$, so the boundary condition $u_{t}(0, t)+\beta(0, t) u_{x}(0, t)=0$ holds in a classical sense.

## Stochastic volatility models

Stock price:

$$
\left\{\begin{array}{l}
d X(t)=\sqrt{Y(t)} X(t) d W \\
d Y(t)=\beta(Y(t)) d t+\sigma(Y(t)) d V
\end{array}\right.
$$

$d W d V=\rho d t$.
We assume that $X$ is absorbed at 0 , and that $\beta(0) \geq 0$,
$\sigma(0)=0$ and linear growth conditions.
Theorem
The function $u(x, y, t)=E_{x, y, t} g(X(T))$ is the unique classical solution to the backward equation

$$
\left\{\begin{array}{l}
u_{t}+\frac{1}{2} x^{2} y u_{x x}+\rho x \sqrt{y} \sigma u_{x y}+\frac{1}{2} \sigma^{2} u_{y y}+\beta u_{y}=0 \\
u(x, y, T)=g(x) \\
u(0, y, t)=g(0) \\
u_{t}(x, 0, t)+\beta(0) u_{y}(x, 0, t)=0 .
\end{array}\right.
$$

## The Kolmogorov (Fokker-Planck) forward equation

Time-homogeneous diffusion:

$$
d X_{t}=\beta\left(X_{t}\right) d t+\sigma\left(X_{t}\right) d W
$$

with $b(0) \geq 0$ and $\sigma(0)=0$
Density:

$$
p(x, y, t):=P_{x}\left(X_{t} \in d y\right) / d y
$$

The density satisfies the backward equation

$$
\left\{\begin{array}{l}
p_{t}=\frac{1}{2} \sigma^{2} p_{x x}+\beta p_{x} \\
p(x, y, 0)=\delta_{y}(x) \\
p_{t}(0, y, t)=\beta(0) p_{x}(0, y, t)
\end{array}\right.
$$

The density also satisfies the forward equation

$$
\left\{\begin{array}{l}
p_{t}=\frac{1}{2}\left(\sigma^{2} p\right)_{y y}-(\beta p)_{y} \\
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What about boundary conditions?

## The CIR-process

$$
d X_{t}=\left(b-a X_{t}\right) d t+\sigma \sqrt{X_{t}} d W
$$

- If $b>\sigma^{2} / 2$, then $\lim _{y \rightarrow 0} p(x, y, t)=0$. Thus $p(x, 0, t)=0$ can be used as boundary condition.


## The CIR-process

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- If $b>\sigma^{2} / 2$, then $\lim _{y \rightarrow 0} p(x, y, t)=0$. Thus $p(x, 0, t)=0$ can be used as boundary condition.
- If $b<\sigma^{2} / 2$, then $\lim _{y \rightarrow 0} p(x, y, t)=\infty$. It is then unclear what boundary conditions to impose.


## A Symmetry Relation

Let $m(x)=\frac{2}{\sigma^{2}(x)} \exp \left\{\int^{x} \frac{2 \beta(z)}{\sigma^{2}(z)} d z\right\}$ be the density of the speed measure.

Theorem
The density satisfies

$$
m(x) p(x, y, t)=m(y) p(y, x, t)
$$

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Note: Using the symmetry relation, we circumvent the difficult issue of boundary conditions for the forward equation. Instead, we solve the backward equation (with boundary conditions specified above), and then use the symmetry relation.

## Boundary asymptotics

As a by-product we also get the exact boundary asymptotics.
Theorem
The density $p(x, y, t)$ satisfies $p(x, y, t) / m(y) \rightarrow C(t)$ as $y \rightarrow 0$ for some $C^{1}$ function $C(t):=p(0, x, t) / m(x) \geq 0$.
(i) If $\beta(0)>0$, then $C(t)$ is positive for $t>0$.
(ii) If $\beta(0)=0$, then the function $C(t) \equiv 0$. In this case, define

$$
D(t)=\lim _{y \downarrow 0} \frac{p(x, y, t)}{y m(y)}
$$

(iia) If there exists a constant $\varepsilon>0$ such that $\sigma(x) \geq \varepsilon x^{1-\varepsilon}$ for $0<x<\varepsilon$, then $D(t)$ is positive for $t>0$.
(iib) If $\beta(0)=0$ and there exists $\varepsilon>0$ such that $\sigma(x) \leq \varepsilon^{-1} x$ for $x \in(0, \varepsilon)$, then $D \equiv 0$.

## Sketch of proof

The above symmetry relation was also proved by Ito-McKean. Their proof is rather sketchy and relies on spectral decomposition of solutions to PDEs. Our proof is more elementary and uses the known regularity in the backward variable and integration by parts.
Fix $x$ and let $\tilde{p}(y, t):=\frac{m(y)}{m(x)} p(y, x, t)$, and let

$$
u(y, t)=E_{y, t} g\left(X_{T}\right)
$$

for some $T>0$ and some function $g$.

Integration by parts:

$$
\begin{aligned}
& \qquad \int_{0}^{\infty} \int_{0}^{T} u_{t}(y, t) \tilde{p}(y, t) d t d y= \\
& \int_{0}^{\infty} u(y, T) \tilde{p}(y, T)-u(y, 0) \tilde{p}(y, 0) d y-\int_{0}^{\infty} \int_{0}^{T} u(y, t) \tilde{p}_{t}(y, t) d t d y . \\
& \text { and }
\end{aligned}
$$

$$
\begin{aligned}
& -\int_{0}^{\infty} \int_{0}^{T} u_{t}(y, t) \tilde{p}(y, t) d t d y \\
= & \int_{0}^{\infty} \int_{0}^{T}\left(\alpha u_{y y}(y, t)+\beta u_{y}(y, t)\right) \tilde{p}(y, t) d t d y \\
= & \int_{0}^{T} \int_{0}^{\infty} u(y, t) \tilde{p}_{t}(y, t) d y d t
\end{aligned}
$$

Thus

$$
\int_{0}^{\infty} u(y, 0) \tilde{p}(y, 0) d y=\int_{0}^{\infty} u(y, T) \tilde{p}(y, T) d y .
$$

Since $\tilde{p}(y, 0)=\delta_{x}(y)$ and $u(y, T)=g(y)$, we have

$$
E_{x, 0} g\left(X_{T}\right)=\int_{0}^{\infty} g(y) \tilde{p}(y, T) d y .
$$

This holds for any function $g$, so $\tilde{p}(y, T)$ is a density of $X_{T}$.

## Example: CIR

$$
\begin{gathered}
d X_{t}=\left(b-a X_{t}\right) d t+\sigma \sqrt{X_{t}} d W \\
m(x)=\frac{2}{\sigma^{2}} x^{2 b / \sigma^{2}-1} \exp \left\{\frac{2 a(1-x)}{\sigma^{2}}\right\}
\end{gathered}
$$

The density satisfies

$$
p(x, y, t)=\frac{m(y)}{m(x)} p(y, x, t) \sim C(x, t) y^{2 b / \sigma^{2}-1}
$$

for small $y$.

## Example: CEV

The Constant Elasticity of Variance model is given by

$$
d X_{t}=\sigma X_{t}^{\gamma} d W
$$

for $\frac{1}{2} \leq \gamma<1$. Now $m(x)=\frac{2}{\sigma^{2}} x^{-2 \gamma}$. The density satisfies

$$
p(x, y, t)=\frac{m(y)}{m(x)} p(y, x, t) \sim C(x, t) y^{1-2 \gamma}
$$

for small $y$.

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## Thank you for your attention!

